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Design and operation of LongBo: a 2 m long drift liquid argon TPC

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ABSTRACT: In this paper, we report on the design and operation of the LongBo time projection chamber in the Liquid Argon Purity Demonstrator cryostat. This chamber features a 2 m long drift distance. We measure the liquid argon purity using cosmic ray muons and confirm the electron drift lifetime is at least 6 ms as measured using purity monitors. LongBo uses cold electronics with discrete circuitry for the preamplifiers, and we measure a signal-to-noise (S/N) ratio of 30 at a drift field of 350 V/cm. An ASIC version of the cold preamplifiers was implemented on 16 readout channels, and we measure a S/N ratio 1.4 times larger for these channels.

KEYWORDS: LArTPC.

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1. Introduction

Liquid argon time projection chambers (LArTPCs) provide excellent spatial and calorimetric resolutions for measuring the properties of neutrino interactions above a few MeV. Conventional liquid argon vessels, such as ICARUS [1], ARGONTUBE [2], *etc.*, are evacuated to remove water, oxygen and nitrogen contaminants present in the ambient air prior to filling with liquid argon. However, as physics requirements dictate larger cryogenic vessels to hold bigger detectors, the mechanical strength required to resist the external pressure of evacuation becomes prohibitively costly. The Liquid Argon Purity Demonstrator (LAPD) [3] was an R&D test stand at Fermilab designed to determine if electron drift lifetimes adequate for large neutrino detectors could be achieved without first evacuating the cryostat. The electron drift lifetime in this test measured with purity monitors was greater than 6 ms without initial evacuation of the cryostat; that lifetime exceeds the value of approximately 2 ms required for future large LArTPCs with drift distances on the scale of a few meters.

After achieving the required electron drift lifetimes, the LAPD cryostat was emptied and a TPC of 2 m drift distance, named LongBo, was installed in the central cryostat region. High liquid argon purity was achieved with the TPC in the cryostat. This paper summarizes the design and operation of the LongBo TPC., ~~which was the first TPC operated in the US to achieve a drift distance of 2 m.~~

2. Construction of the LongBo TPC

The LongBo TPC has a cylindrical sensitive volume with a 25 cm diameter and a 2 m length. LongBo extended the drift distance of the "Bo" TPC that was made a few years earlier with a 50 cm length [4]. LongBo has three sense wire planes, with wires oriented at 60 degrees from each other. There are 48 instrumented wires in each plane and the wires are 4.7 mm apart and made of 125 micron diameter beryllium copper. They were hand-soldered to a 3 mm thick copper-clad G10 board for each plane. The copper cladding was routed in a pad pattern on a Fermilab CNC router. A set of individual wires (shown with white insulation in figure 1) were soldered to the readout boards and terminated in a connector matched to a corresponding one on the preamplifier motherboards. The preamplifiers, often referred to as "cold preamplifiers" since they operated in the 87 K liquid argon, were placed around the TPC just outside the field cage. ~~; an arrangement dictated by the size of the cryostat flange on which the TPC was mounted.~~ Olefin insulated flat cables took the signals from the motherboards to the signal feedthrough. The planes were biased from external voltage sources via a feed through. The cathode was made from copper mesh, chosen to allow free argon flow along the TPC axis. The mesh was stretched inside a ring made of 19 mm diameter 3 mm thick stainless steel tube, which supports the TPC. A stainless steel tube elbow was welded to this ring and served to connect to the high voltage feed through, described in Section 3, via a cup and spring-finger connection.

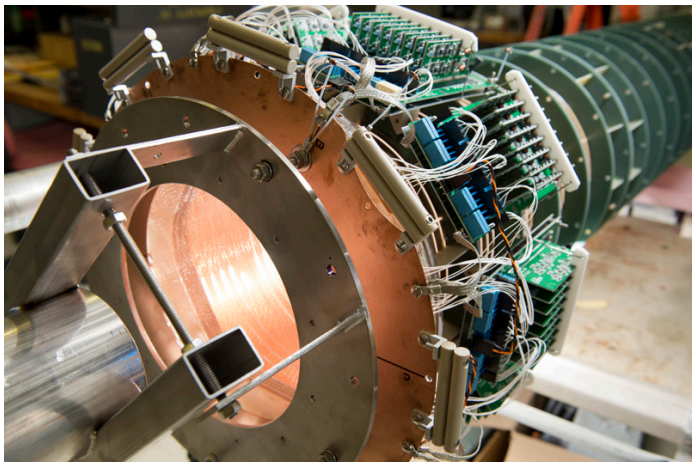


Figure 1: The wire planes and cold preamplifiers.

A uniform electric drift field was provided to the Bo TPC by a cylinder made of a copper-clad, 0.6 mm thick, G10 sheet, which is 50 cm long by 78 cm wide. The sheet was routed to create 18 mm wide copper strips, separated by 1.5 mm. The sheet was rolled into a 25 cm diameter tube which was inserted into a set of 7 rings routed from 4 mm thick G10. Two strings of 100 MOhm voltage divider resistors, Ohmite model 104, rated at 10 kV, 1 W, were soldered to the inner surface of the copper sheet. Only one string is necessary to create the electric field; the second is for redundancy. Figure 2 shows the resistors soldered to the inner surface of the copper sheet.

Bo was converted to LongBo by inserting three additional drift field sections, each 50 cm long, for a total drift length of 2 m. The drift field sections were connected together using threaded rod

Figure 2: The resistors soldered to the inner surface of the copper sheet.

to make the full 2 m drift cage. Beryllium copper spring stock were soldered to the ends of each cylinder to assure inter-connection at the joints. Figure 3 shows the assembled TPC with high voltage feedthrough and electronics.

3. The High Voltage System

To provide the electric field to drift the signal electrons, a high voltage was applied to the cathode mesh on the bottom of the TPC. The high voltage was generated outside of the cryostat by a Glassman LX150N12 power supply [5]. The supply was controlled remotely by a program and the normal operation current is much less than 1 mA. The supply was set to trip if more than 1.1 mA of current was drawn in order to protect the equipment and personnel who operate the system. Before entering the cryostat, the voltage was passed through a filter pot. This low-pass filter was a sealed aluminum vessel with cable receptacles that electrically connected to eight 10 MΩ resistors submerged in Diala oil. The capacitance for the filter was supplied by the cable to the feedthrough. The purpose of the pot was twofold: it reduced the high-frequency ripple from the power supply, and it limited the energy to the TPC in the case of a high voltage discharge.

A feedthrough transmits the high voltage into the cryostat and to the receptacle cup of the cathode plane. This feedthrough is shown in figure 4 and was modeled after the feedthrough used in the ICARUS experiment [1] with a stainless steel inner conductor insulated radially by ultra high molecular weight polyethylene (UHMW PE), and surrounded by a stainless steel ground tube. Grooves were added to the exposed UHMW PE to reduce surface currents. Initially, conducting shielding cups, shown in figure 4b, were installed to reduce the field along the feedthrough. Midway through running, the cups were removed to determine what effect they had on performance. No change was seen. To ensure good electrical contact, the feedthrough has a spring tip that is inserted into a receptacle cup attached to the cathode plane.

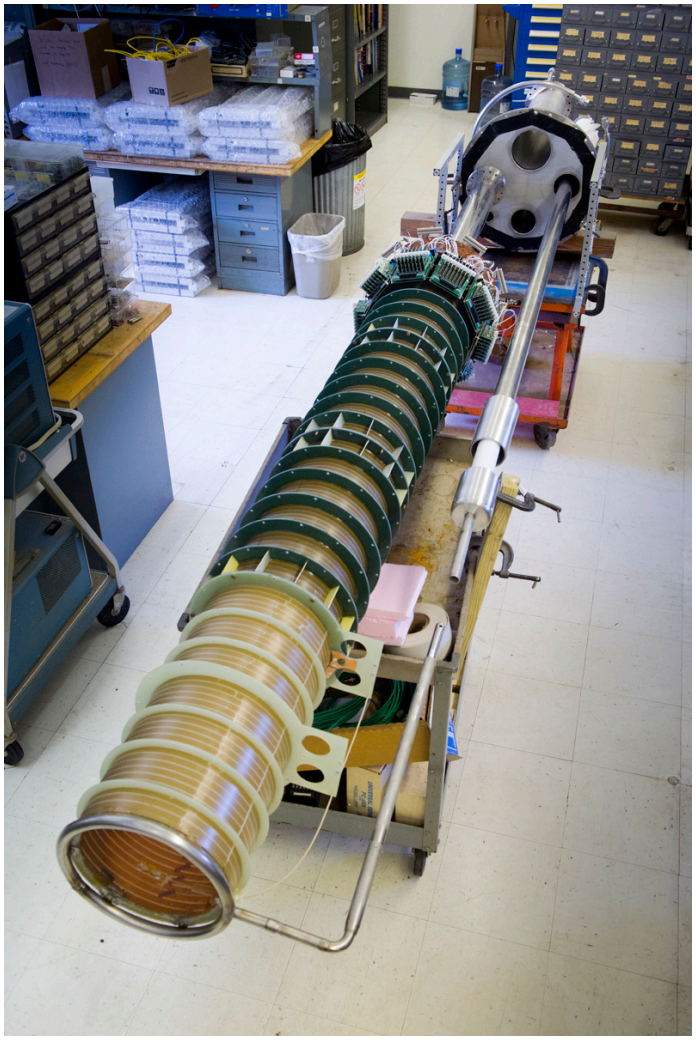


Figure 3: Assembled LongBo TPC.

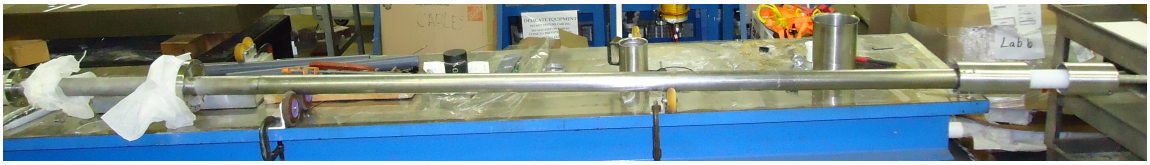
4. Electronics

4.1 Cold Electronics

The LArTPC technology was pioneered by the ICARUS collaboration using warm electronics with a dual-jFET charge integrating front-end preamp. For each wire of the detector the feed-through passed DC bias-voltage into the cryostat and passed the TPC signals back to the electronics. Early development in the US used a similar warm preamplifier on a 3-plane LArTPC detector, Bo [4], built by Fermilab and used a 32-channel ADF2 digitizer built by MSU for the D0 experiment. At the same time the ArgoNeuT [6] LArTPC was built and used the same warm preamplifier design readout.

~~The readout electronics used in LongBo are the result of earlier development efforts by MSU for the Bo [4] and ArgoNeuT [6] TPCs. These systems employed preamplifiers positioned outside of the cryostat.~~ The development of cold electronics for a LArTPC was driven by fact that at the

(a)



(b)

Figure 4: Photographs of the HV feedthrough without (a) and with (b) the electric field shielding cups.

87 K temperature of liquid argon there is less intrinsic noise in a MOSFET preamp than in the best warm preamp. Additionally, by mounting the preamp on the TPC the noise generated by a cable capacitance can be eliminated. A cold MOSFET preamp-filter card was designed and ~~144 channels were~~ built for the Bo TPC to replace the original warm preamplifiers. In this first application of a MOSFET front-end on a LArTPC, the preamp-filter was implemented with a 2.1 μs peaking time. These "discrete" MOSFET preamp-filter cards were used on the LongBo LArTPC described in this paper.

A MOSFET ASIC preamp-filter for use at cryogenic temperatures was designed by BNL for LBNE [10]. It is also the readout of the TPC detector for the MicroBooNE experiment [11]. A version (V4) of this 16-channel ASIC became available in 2012 and was implemented for the LongBo TPC. One 16-channel preamp-filter card was modified to use a section of the MicroBooNE motherboard that serviced a single ASIC as a mezzanine board. The input/output connectors, bias-voltage distribution and decoupling capacitors of the card were retained. An ASIC control box was built which set the gain and peaking time parameters of the ASIC. The ASIC modified preamp-filter card was installed in place of one induction plane card with the signals processed by the same digitizing electronics.

4.2 Signal to Noise Ratio

The signal to noise ratio is an important metric of performance for a liquid argon detector. The absolute numerical value depends upon wire spacing, wire capacitance, readout electronics and drift field strength. Because 16 channels of the LongBo readout electronics were instrumented with the BNL ASIC preamplifier, a direct performance comparison can be made between the ASIC version of the preamplifier and the discrete CMOS version used for the remaining channels.

To measure the noise for each channel, a special run of 200 triggers was taken with a random trigger and no drift field. For each channel, the measured ADC counts for all 200 triggers were combined and the distribution fit to a Gaussian function. The width of the best fit Gaussian is reported as the noise in figure 5. The measured noise in Fig. 5 can also be determined from regular

data runs with an external muon trigger and nominal drift field if signals from cosmic ray muons are removed from the acquired waveforms.

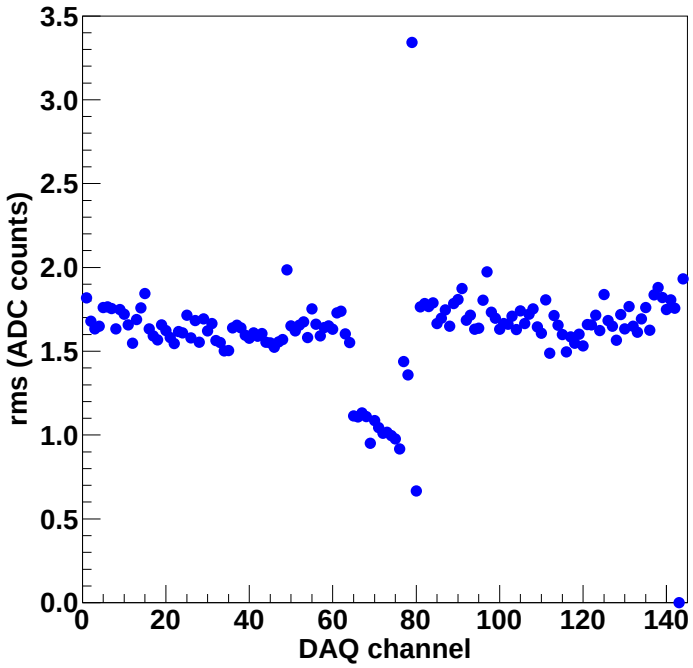


Figure 5: Measured noise as a function of DAQ channel number. Channels 1-48 are the first induction plane, channels 49-96 are the second induction plane and channels 97-144 are the collection plane. Channels 65-80 are instrumented with the BNL ASIC preamplifier, all other channels are instrumented with the MSU discrete CMOS preamplifier. After the first HV breakdown, four channels of the ASIC readout (channels 77-80) exhibited more noise than before and it was no longer possible to change their parameter settings. All of the MSU discrete CMOS channels survived the HV breakdowns.

The signal strength was extracted from straight muon tracks acquired with the cosmic ray muon trigger. The signal to noise ratio of the discrete electronics chain was determined from signals on the collection plane wires. Clean tracks were selected from the first 200 triggers in each run considered using the event display. Only tracks whose projection onto the wire plane under study was nearly perpendicular to wire direction and centered on the wire plane were selected.

The signal strength is recorded as the maximum ADC count in the waveform. The outermost two wires on either edge of the wire plane were excluded from the analysis as a fiducial volume cut intended to reject hits close to the field cage where the drift field may not be uniform and the pulse heights are noticeably reduced. In addition, wires exhibiting pulse heights greater than 1.5 times the average pulse height for the track were interpreted as delta rays and excluded. Finally, wires with more than one hit for a given trigger were excluded.

In the absence of diffusion, the length of the track segment that contributes to the signal on a particular collection plane wire is $\frac{\Delta w}{\sin(\theta)\cos(\phi)}$, where $\Delta w = 4.7$ mm is the wire spacing and θ and ϕ are the reconstructed track angles defined in § 6.1. The individual measured signal strengths were reduced by a factor of $\sin(\theta)\cos(\phi)$ before taking the average of all *good* collection wire values from the 10-12 event sample to obtain the value of the signal strength $S = 50.3$ ADC counts and a signal to noise ratio of $50.3/1.69 \sim 30$ for the discrete CMOS electronics chain.

The signal to noise ratio provides a good comparison of the performance of the different amplifier flavors. The ASIC preamplifier is connected to the center 16 wires of the middle (induction) wire plane, so induction plane signals must be compared. For these bipolar signals, the signal strength used is the difference between the pedestal and the minimum ADC count of the waveform, or the depth of the dip in the bipolar signal shape. Tracks whose projection onto the middle wire plane is perpendicular to the wire direction were selected ($\phi \sim 60^\circ$) and the extracted signal strengths were normalized by the factor $\sin(\theta)\cos(\phi - 60^\circ)$. The ratio of the average signal strength for ASIC and discrete channels is reported in table 1 for different values of the shaping time constant. Both the measured noise and the signal to noise ratio increase with the shaping time constant.

Table 1: Comparison of signal and noise for discrete CMOS preamplifier and the ASIC preamplifier for different settings of the shaping time constant of the ASIC. The discrete CMOS preamplifiers have a fixed shaping time of 2 μs. The gain setting of the ASIC was 25 mV/fC for all measurements tabulated.

Shaping time (μs)	S	$S_{discrete}$	N	$N_{discrete}$	$(S/N)/(S/N)_{discrete}$
0.5	10.1	36.1	0.55	1.69	0.9
1.0	19.3	35.1	0.75	1.69	1.2
2.0	33.6	35.5	1.10	1.69	1.4
3.0	40.3	35.8	1.35	1.69	1.4

Figure 6 shows the shape of the amplified signal on a single wire for different shaping times, with the the signals scaled to have the same peak height to facilitate comparison. There is little if any difference in the pulse shapes - the shape is dominated by the width of the pulses before amplification. Changing the shaping time constant changes the overall gain of the amplifier.

4.3 Trigger

The presence of a through-going muon was detected by three sets of scintillation counters placed around the exterior of the LAPD tank. Each set consisted of two groups of counters placed on opposite sides of the tank, with the line intersecting the groups of counters passing roughly perpendicular to the wire direction for one wire plane. There were 8 counters (4 pairs) in each of the

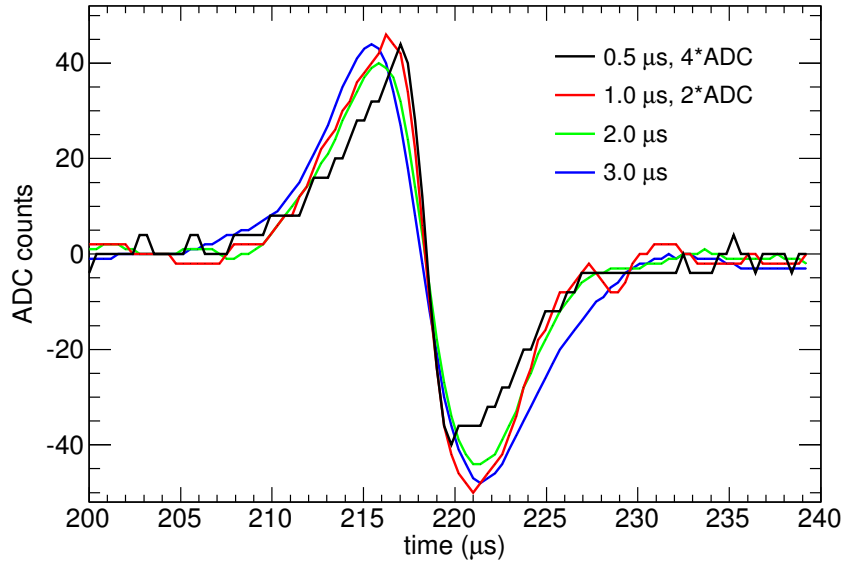


Figure 6: Raw signal shapes for induction plane signals with different settings of the ASIC shaping time constant are shown. The horizontal axis has an arbitrary zero so that the pulses are shown with the zero crossing point artifically aligned. The tracks from which these signals are taken form an angle of roughly 20 degrees with the plane of the wires, or $\theta \sim 70$. The angle between the projection of the track onto the wire plane and the wires themselves is ~ 80 degrees.

6 groups, stacked in a column for an effective area of scintillator roughly 3 m by 0.4 m, for a total of 48 counters. One control and concentrator unit (CCU) controls 48 counters setting the gain of the photomultiplier tube (PMT) and a discriminator threshold [12]. A coincidence was required between counters on opposite side of the tanks, and the track angle coverage is shown in § 6.1.

5. Operation

The LAPD cryostat is an industrial low pressure storage tank. The cryostat has a flat bottom, cylindrical sides, and a dished head. The cryostat diameter is 3.0 m and the cylindrical walls have a 3.0 m height. Figure 7 shows the positions of the TPC and the high voltage feed through inside the LAPD cryostat. After purging the cryostat with argon gas and then filtering the argon gas for several days, as described in [3], the cryostat was filled with liquid argon. The high voltage was raised after the electron drift lifetime passed 0.3 ms. The TPC data taking lasted for 8 months.

An issue we encountered in the LongBo run was that the high voltage was unstable. The cause of the instability is unknown, but was present during the entire data-taking period. This instability forced the running of the TPC at a lower high voltage than desired, ultimately reaching a maximum -75 kV, with occasional trips, instead of the desired 100 kV for a 2 m drift TPC. Even with the lower-than-desired high voltage, several hundred thousand clear cosmic ray muon tracks were recorded. The data were used to study the electronics performance and liquid argon purity.

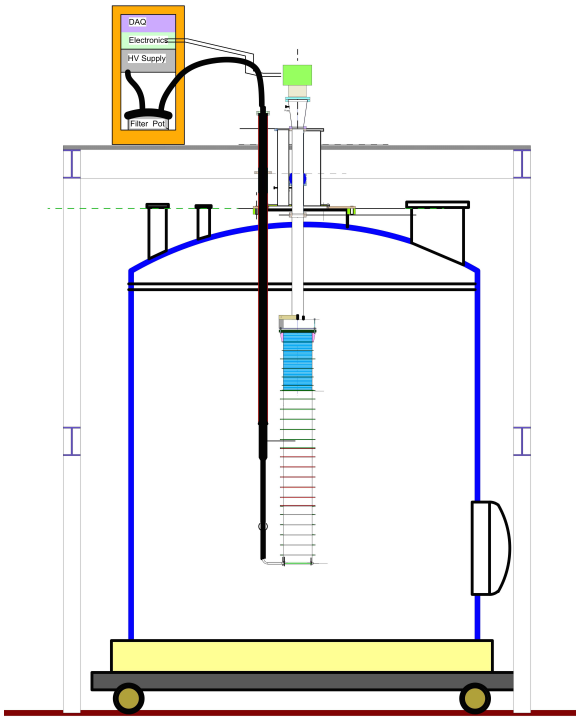


Figure 7: The TPC and high voltage feed through inside the LAPD cryostat.

6. Results

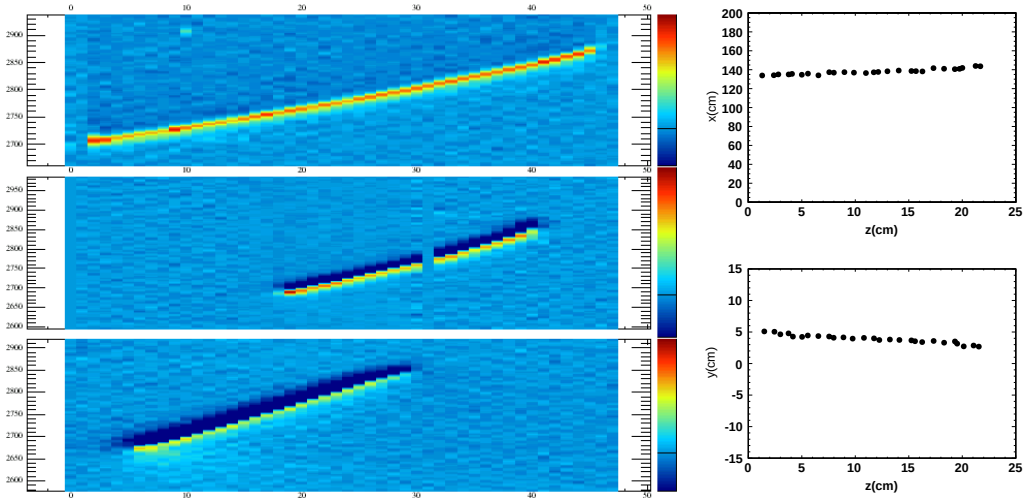
6.1 Measurement of Electron Attenuation Using Cosmic Ray Muons

When cosmic ray muons pass through the liquid argon, they deposit energy through ionization and scintillation. The ionization electrons are drifted by the electric field and collected by the wire planes. The variation of the energy deposition along the muon track is small for the energetic muons. By examining the signal recorded by each wire as a function of electron drift time, one can measure the attenuation of ionization electrons along the drift distance and determine the electron drift lifetime. This method was used by ArgoNeuT to derive an electron drift lifetime [6].

We analyzed cosmic ray muon data taken during one full cycle of LAPD running, from when the liquid argon started recirculating through the filters to when it stopped. The high voltage we applied to the cathode was -70 kV during this period, which produced an electric field of 350 V/cm in the TPC volume. In this section, we present results on the measurements of electron attenuation using these data.

There are 491 966 triggered events during this run period. We use the LARSOFT software package [13] to reconstruct cosmic ray muon events. The automated reconstruction first converts the raw signal from each wire to unipolar pulse by removing electronics and electric field responses through deconvolution, and then finds hits by fitting a Gaussian to the resulting pulse. The hits from each plane of the TPC are grouped into clusters in drift time and wire number. Three dimensional tracks are constructed from pairs of line-like clusters in each plane. There are 274 491 events with

at least one reconstructed track. Figure 8 shows one example event where a muon track is visible.



(a) Raw data (b) Reconstructed track

Figure 8: Run 293 Event 2225. (a) Drift time versus wire ID for raw wire signal from 3 wire planes. Red color represents large positive charge and blue color represents large negative charge. (b) Reconstructed track points in the x-z and y-z projections.

In this analysis we only consider the 223 194 events with exactly one reconstructed track that has at least 5 reconstructed 3-dimensional points along the track trajectory. We do not use events with multiple muon tracks because we cannot determine the track start time, t_0, for each track. Figure 9 shows the angular distributions of the reconstructed tracks. The track angle is determined by the reconstructed track start and end points, assuming the track is a straight line. The zenith angle θ is peaked around 60° and the azimuthal angle ϕ is determined by the trigger counter locations. Figure 10 shows the reconstructed track start point x, in the drift direction, versus the θ angle; $x = 0$ is near the wires. The muons normally enter the TPC in the middle, but they can be close to the wire planes and cathode.

The determination of the argon purity is performed using the collection plane hits on the track. We first select events with well reconstructed and clean tracks, where

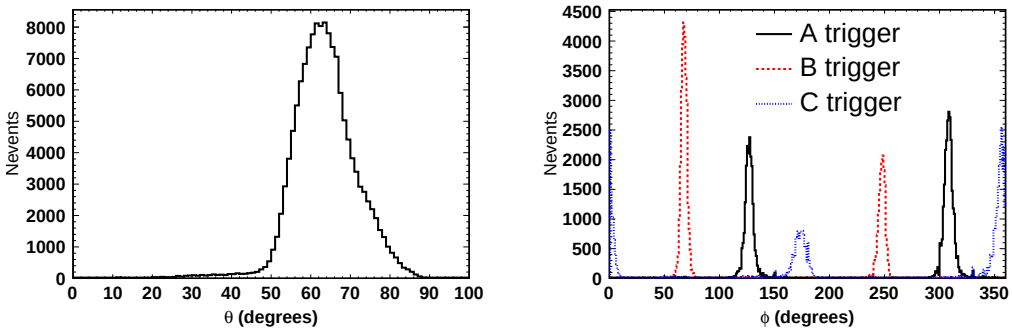
1. the RMS distance of reconstructed 3D points with respect to a straight line determined by track start and end points should be less than 6.3 cm;

2. the number of hits that are not associated with the reconstructed track should be less than 11;

3. track length is required to be greater than 15 cm;

4. track length $\times \cos\phi$, the projection along the collection plane trigger direction, is required to be greater than 10 cm;

5. θ is required to be tween 50° and 70°.



(a) Track θ (b) Track ϕ

Figure 9: Reconstructed track angles: (a) θ - with respect to the vertical direction, while $\theta = 0$ is at the zenith; (b) ϕ, events taken with different triggers have distinct ϕ distributions.

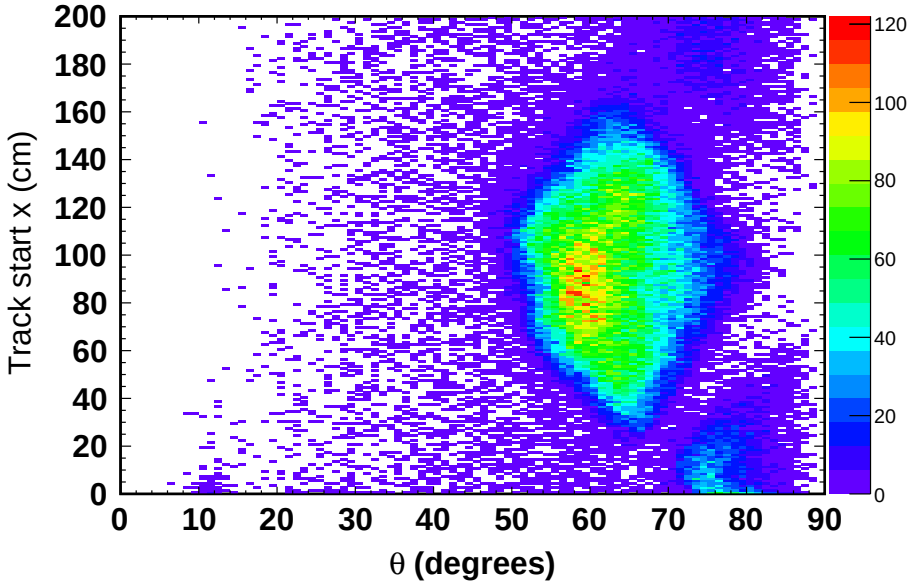


Figure 10: Reconstructed track start point in the drift direction (x) versus track angle (θ). The muons normally enter the TPC in the middle, but they can be close to the wire planes and cathode.

The selection was chosen to maximize the good track fraction. We then select all hits on the collection plane associated with the tracks. We require only one hit on each wire. If there is one hit on each of three contiguous wires that has $dQ/dx > 2000$ ADC/cm, we do not use those three hits. This requirement is meant to remove delta rays.

For each selected hit, we look at raw wire signal for that hit in the region $[t - 3 \times \sigma, t + 3 \times \sigma]$, where t is the reconstructed hit time and σ is the reconstructed hit width. We first find the peak of the raw digits in ADC, we then sum all the digits above a threshold of 10% of the peak to get the charge of the hit. This threshold is chosen based on measured signal to noise ratio. By using

raw wire signal we remove uncertainties introduced by signal shaping in the hit reconstruction. We then divide the hit charge by the track pitch, which is defined as the wire pitch over dot product of the track direction and the direction normal to the wire direction in the wire plane, to get dQ/dx for the hit.

Figure 11 shows the hit dQ/dx distributions when the argon purity was poor (left) and good (right). The top plots show dQ/dx versus electron drift time distributions. The middle plots show the dQ/dx distributions for hits with drift time betwen 456 and 507 μs, indicated as gray bands in the top plots. We fit a Landau convoluted with Gaussian function to the dQ/dx distribution. The bottom plots show the most probable value (MPV) from the Landau fit as a function of drift time. The signal decreases as drift time increases as expected. We fit an exponential function to the data points:

$$dQ/dx = e^{-at_d+b}, \tag{6.1}$$

where t_d is the drift time and a is the attenuation constant. Figure 12 shows $dQ/dx_0 = dQ/dx(t = 0) = e^b$ as a function of time and the fluctuation is small.

LAPD reported the measurement of electron attenuation using purity monitors [3]. The result was presented as Q_A/Q_C as a function of time, where Q_A/Q_C represents the fraction of electrons generated at the purity monitor cathode that arrive at the anode. It is determined by the electron drift time (t_d) and electron lifetime (τ) or attenuation constant (a):

$$Q_A/Q_C = e^{-t_d/\tau} = e^{-at_d}. \tag{6.2}$$

In order to compare the measurement using TPC data with the measurement using purity monitors, we calculate the equivalent Q_A/Q_C using the attenuation constant measured by the TPC data and $t_d = 0.31$ ms, which is the electron drift time from the cathode grid to the anode grid in the purity monitor. Figure 13 shows the comparison between Q_A/Q_C measured by the purity monitors and the calculated Q_A/Q_C using attenuation constant measured using the TPC data. The statistical uncertainties from the original Landau fits are propagated through the attenuation fits to figure 13, which are negligible compared with the systematic errors discussed below. The sources of systematic error on Q_A/Q_C measured by the purity monitor are discussed in Ref. [3] and shown on the plot as a combined 5% absolute uncertainty.

In order to evaluate the systematic uncertainties on the purity measurement using TPC data, we test the same method using a 100 000 muon track sample in the LongBo TPC generated with the energy spectrum and the angular distribution of cosmic rays by LARSOFT. The simulation includes the muon ionization and recombination effect, the free electron longitudinal and transverse diffusion, the signal attenuation due to the impurities in liquid argon and the electronic response. The method has been tested for electron lifetime in the 0.5 - 30 ms interval. The results, summarized in figure 14, show the measurement is reliable in a large range of liquid argon purity values. We fit a straight line to the points and we add an error of 0.002 to the points in addition to the statistical errors in order to get a reasonable χ^2/ndf value. The linear fit result

$$(Q_A/Q_C)^{Measured} = -0.006 + 1.012 \times (Q_A/Q_C)^{True} \tag{6.3}$$

is used to correct the data measurements to obtain the results shown in figure 13. We attribute 0.002 as the systematic error from the method to derive the Q_A/Q_C measurement.

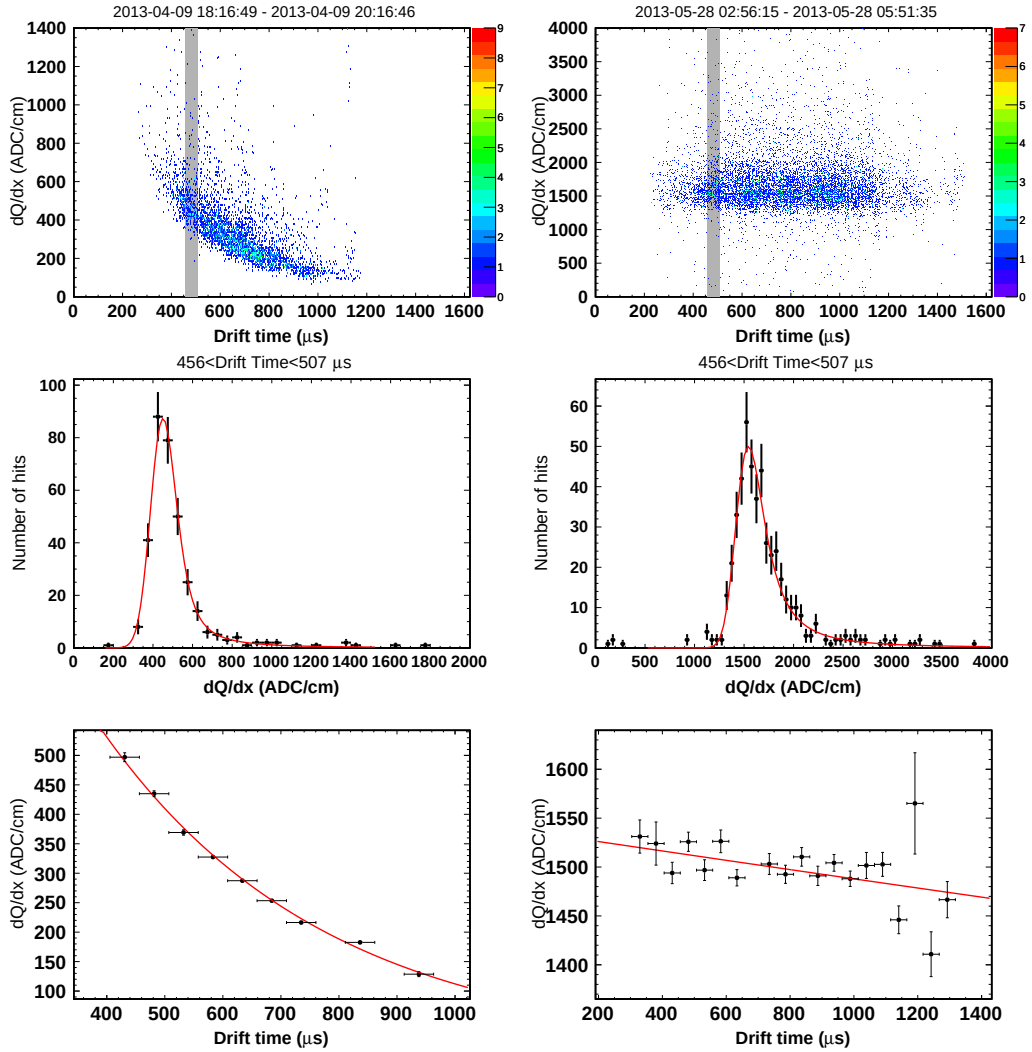


Figure 11: dQ/dx distributions of hits using data taken in two periods of time when the electron lifetime was 0.4 ms (left) and 30 ms (right), respectively. Top: Scatter plots of dQ/dx as a function of electron drift time; Middle: dQ/dx distributions for hits with drift time between 456 and 482 μs; Bottom: dQ/dx from Landau fit as a function of drift time.

Figure 15 shows the data Q_A/Q_C distribution when liquid argon purity is good. We fit a Gaussian to the data distribution. The sigma of the Gaussian fit indicates the system instability and is take as one source of systematic error as well as the error on the fitted Gaussian mean value. The total systematic error on the measured Q_A/Q_C is

$$\sigma_{\text{sys}} = 0.002_{\text{method}} \oplus 0.009_{\text{system}} \oplus 0.001_{\text{error on the mean}} = 0.0096, \tag{6.4}$$

and is shown in figure 13.

We use the Bayesian method [14] to determine the lower limit on $s = Q_A/Q_C$. The likelihood

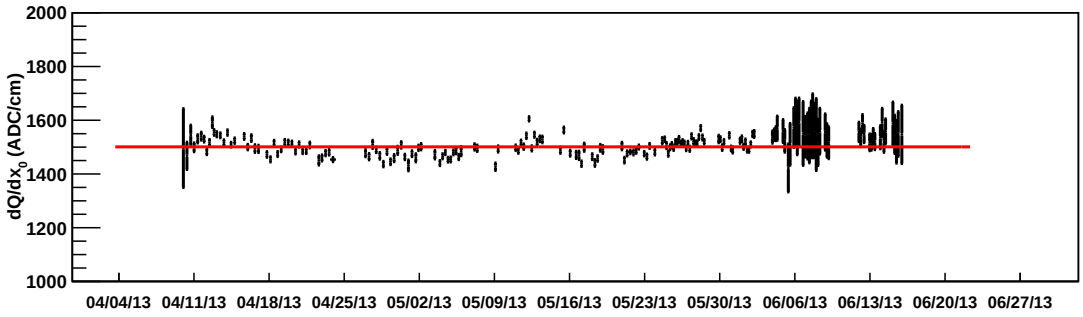


Figure 12: dQ/dx_0 as a function of time.

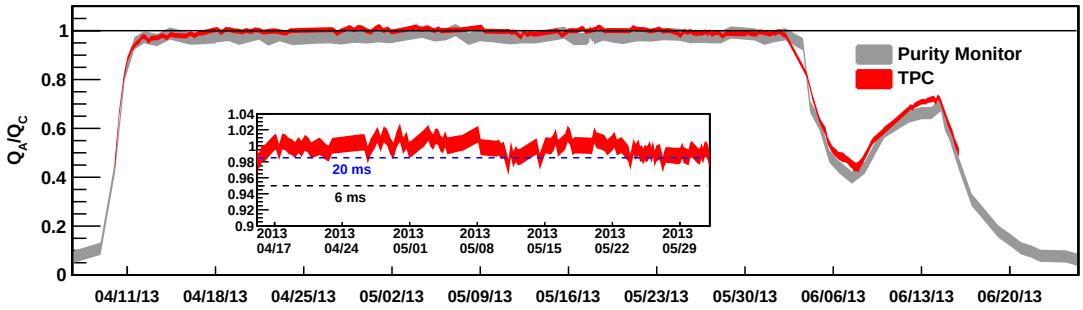


Figure 13: Comparison between Q_A/Q_C measured by the purity monitors and the calculated Q_A/Q_C using attenuation constant measured using the TPC data. The details of the TPC results are shown in the inset compared with the values $Q_A/Q_C = 0.95$ and 0.985 which corresponds to electron drift lifetime $\tau = 6$ and 30 ms, respectively.

for s is given by the Gaussian distribution with mean 0.9974 and sigma 0.0096:

$$P(s) = \frac{1}{\sqrt{2\pi}0.0096} e^{-\frac{(s-0.9974)^2}{2\times 0.0096^2}}. \tag{6.5}$$

For s we use the prior:

$$\pi(s) = \begin{cases} 1 & s \leq 1 \\ 0 & s > 1 \end{cases} \tag{6.6}$$

to take into account the physics boundary of Q_A/Q_C. A lower limit s_{lo} at confidence level $1-\alpha$ can be obtained by requiring

$$\alpha = \frac{\int_{-\infty}^{s_{lo}} P(s)\pi(s)ds}{\int_{-\infty}^{+\infty} P(s)\pi(s)ds} = \frac{\int_{-\infty}^{s_{lo}} P(s)ds}{\int_{-\infty}^{1} P(s)ds}. \tag{6.7}$$

For 95% confidence level ($\alpha = 0.05$) we obtain a lower limit on Q_A/Q_C: $s_{lo} = 0.985$, which corresponds to an electron drift lifetime of 20 ms.

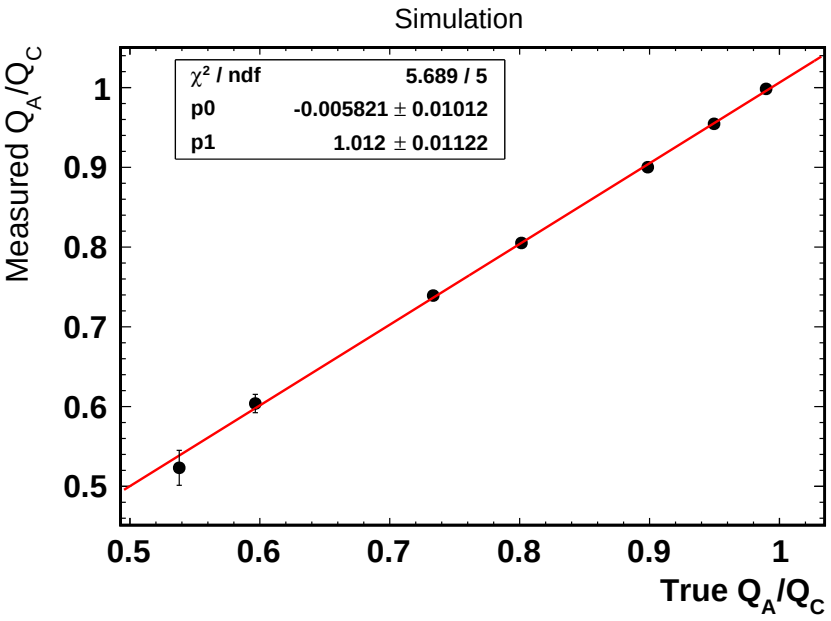


Figure 14: The ratio Q_A/Q_C as measured with the Monte Carlo events as a function of the simulated Q_A/Q_C values.

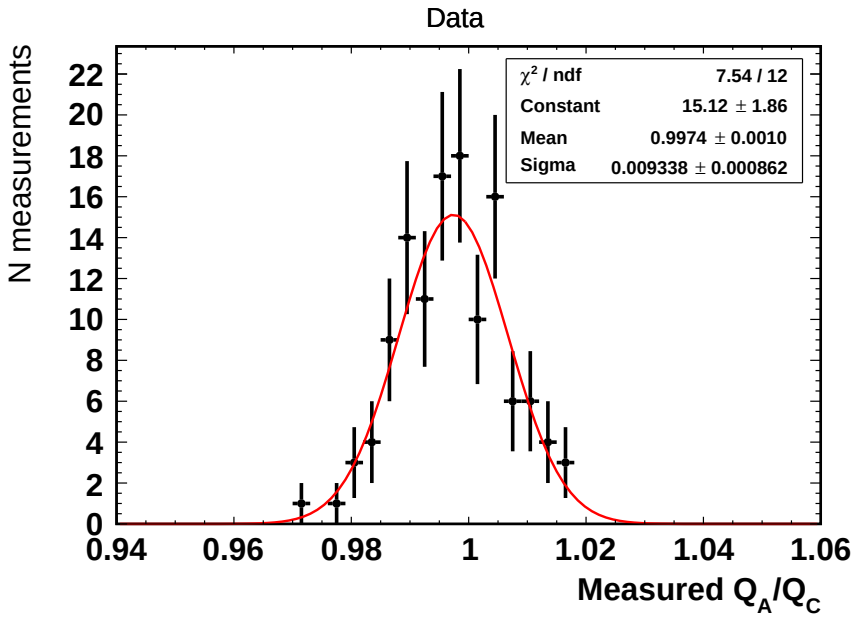


Figure 15: The data Q_A/Q_C distribution when liquid argon purity is good (data points in inset of figure 13).

6.2 High Voltage Stability

The TPC's resistive chain was likely damaged before data taking began. The power supply was

set to turn off if it drew an over-current much greater than twice the operating current of the TPC. During data taking, the cathode high voltage never sustained greater than 80 kV for a significant period with most data being collected with 60 or 70 kV on the cathode. A correlation was noted between greater high voltage stability and greater electronegative contamination in the argon. This trend can be seen in figure 16 where the number of high voltage trips is displayed along with a purity monitor reading as a function of time. While a precise characterization of the dependence is beyond the scope of this paper, the behavior agrees with other studies in the literature that have shown a trend between an increase in electronegative contamination and an increase in liquid argon's dielectric strength [7, 8, 9].

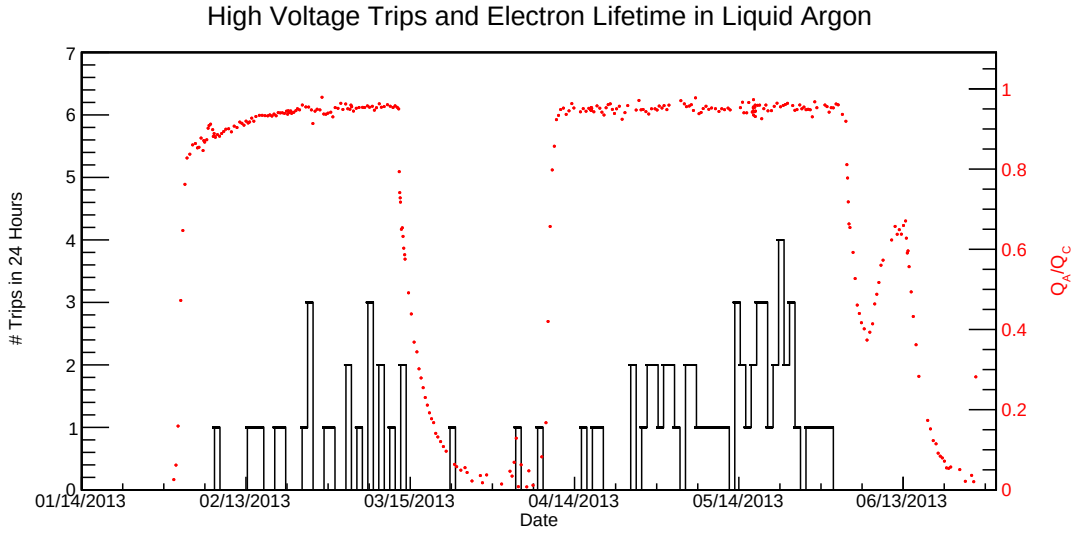


Figure 16: The number of cathode high voltage trips in a 24 hour period and purity monitor values during data taking. The trip counts are given by the black solid line and scale on the left. Trips related to testing the high voltage system were not included in the trip count. A quantification of the argon purity is given by the red dots and the scale on the right.

7. Conclusions

In this paper we have discussed the design and operation of the LongBo TPC in the LAPD cryostat. A ratio $S/N = 50.5/1.69 \sim 30$ was measured for the wires in the collection plane with discrete CMOS electronics. For these data the TPC was operated with a drift field of 350 V/cm. The measured S/N ratio for the ASIC channels was 1.44 times larger than that measured for the discrete channels for the ASIC shaping time of 2.0 μs and a gain setting of 25 mV/fC. The liquid argon purity measurement using cosmic ray muons in the TPC shows the electron drift lifetime is at least 20 ms at 95% confidence level for a long period of time. We also observe a correlation between the voltage stability and the liquid argon purity.

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